

# Summary of Special Relativity in Minkowski Notation (Lecture 11)

Contravariant (upper index)

$$x^\mu = (x^0, x^1, x^2, x^3) = (ct, x, y, z)$$

Covariant (lower index)

$$x_\mu = (x_1, x_2, x_3, x_4) = (ct, -x, -y, -z)$$

Lorentz Transformation

$$x'^0 = \gamma(x^0 - \beta x^1)$$

$$x'^1 = \gamma(-\beta x^0 + x^1)$$

$$x'^2 = x^2$$

$$x'^3 = x^3$$

$$x'^\mu = \Lambda^\mu_\nu x^\nu$$

$$\Lambda^\mu_\nu = \begin{bmatrix} \gamma & -\gamma\beta & 0 & 0 \\ -\gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

$$x'_0 = \gamma(x_0 + \beta x_1)$$

$$x'_1 = \gamma(\beta x_0 + x_1)$$

$$x'_2 = x_2$$

$$x'_3 = x_3$$

$$x'_\mu = (\Lambda^{-1})^\nu_\mu x_\nu$$

$$(\Lambda^{-1})^\nu_\mu = \begin{bmatrix} \gamma & \gamma\beta & 0 & 0 \\ \gamma\beta & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

## Exercise 1

By considering the Lorentz transformations in both directions  $x \rightarrow x'$  and  $x' \rightarrow x$ , for contra- and co-variant  $x$  we

can show that:

$$\left(\Lambda^{-1}\right)^\mu{}_\nu = \frac{\partial x'_\nu}{\partial x_\mu} = \frac{\partial x^\mu}{\partial x'^\nu}$$

$$\Lambda^\mu{}_\nu = \frac{\partial x'^\mu}{\partial x^\nu} = \frac{\partial x_\nu}{\partial x'_\mu}$$

Now we can show that

$$\Lambda^\mu{}_\nu \left(\Lambda^{-1}\right)^\nu{}_\rho = \delta^\mu{}_\rho$$

## Metric Tensor

$$g_{\mu\nu} = g^{\mu\nu} = \begin{bmatrix} 1 & & & \\ & -1 & & \\ & & -1 & \\ & & & -1 \end{bmatrix}$$

$$g_{\mu\lambda} g^{\lambda\nu} = \delta_{\mu}^{\nu} = \begin{bmatrix} 1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{bmatrix}_{4 \times 4}$$

$$x^{\mu} = g^{\mu\nu} x_{\nu} \quad \text{"raises Lorentz index"}$$

$$x_{\mu} = g_{\mu\nu} x^{\nu} \quad \text{"lowers Lorentz index"}$$

## Scalar Product

$$a_{\mu} b^{\mu} = a^{\mu} b_{\mu} = a_{\mu} g^{\mu\nu} b_{\nu} = a^{\mu} g_{\mu\nu} b^{\nu} = a^0 b^0 - a^1 b^1 - a^2 b^2 - a^3 b^3$$

is a Lorentz scalar (invariant)  $a'_{\mu} b'^{\mu} = a_{\mu} b^{\mu}$

## Magnitude

$$x_{\mu} x^{\mu} = (ct)^2 - x^2 - y^2 - z^2 \quad \text{is also Lorentz invariant.}$$

# Summary of Lecture 11    The operator $\partial^\mu$

We define:

$$\partial_\mu = \frac{\partial}{\partial x^\mu} = \left( \frac{1}{c} \frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right) \quad \therefore \quad \partial'_\mu = \frac{\partial}{\partial x'^\mu} = \left( \Lambda^{-1} \right)^\nu{}_\mu \frac{\partial}{\partial x^\nu}$$

$$\partial^\mu = \frac{\partial}{\partial x_\mu} = \left( \frac{1}{c} \frac{\partial}{\partial t}, -\frac{\partial}{\partial x}, -\frac{\partial}{\partial y}, -\frac{\partial}{\partial z} \right) \quad \therefore \quad \partial'^\mu = \frac{\partial}{\partial x'_\mu} = \Lambda^\mu{}_\nu \frac{\partial}{\partial x^\nu}$$

That is:

Differential w.r.t. contravariant,  $x^\mu$ , transforms as covariant 4-vector  
" " covariant,  $x_\mu$ , " " contravariant "