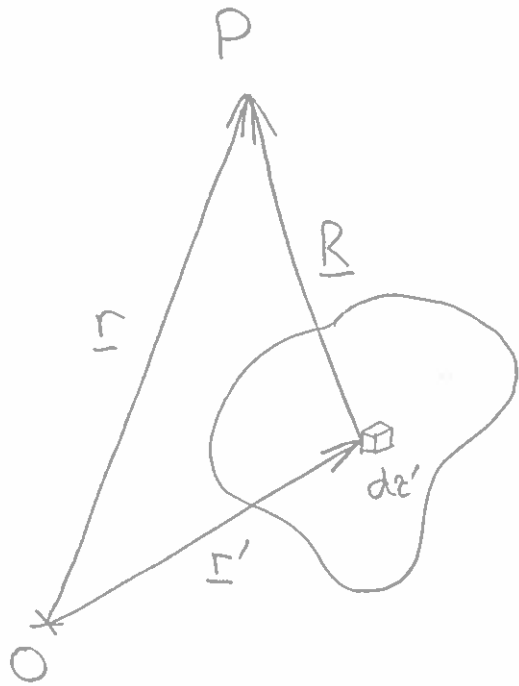


## Summary from Lectures 1 and 2

In Electrostatics we have:



$$V(\underline{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\underline{r}')}{R} d\tau' \quad \text{where } \underline{R} = \underline{r} - \underline{r}'$$

$$\text{We showed that } \nabla\left(\frac{1}{R}\right) = -\frac{\hat{\underline{R}}}{R^2}$$

Therefore, defining  $\underline{E} = -\nabla V$  gave

$$\underline{E}(\underline{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\hat{\underline{R}}}{R^2} \rho(\underline{r}') d\tau'$$

We showed that  $\oint \underline{E} \cdot \underline{da} = \frac{1}{\epsilon_0} \int \rho(\underline{r}') d\tau' = \frac{Q_{\text{enclosed}}}{\epsilon_0}$ ,

or in differential form  $\nabla \cdot \underline{E} = \frac{\rho}{\epsilon_0}$  or  $\nabla^2 V = -\frac{\rho}{\epsilon_0}$

## Reminders

$\underline{r}'$  position of "source" (charge, current)

$\underline{r}$  position at which potential/field evaluated (relative to origin)

$\underline{R}$  position at which potential/field evaluated (relative to "source")

$$\nabla = \left( \hat{x} \frac{\partial}{\partial x} + \hat{y} \frac{\partial}{\partial y} + \hat{z} \frac{\partial}{\partial z} \right)$$

· represents partial differentiation w.r.t. the unprimed coordinates

∴  $\nabla$  operating on any scalar  $f(\underline{r}')$  or vector  $\underline{v}(\underline{r}')$  that is a function only of the primed coordinates gives zero.