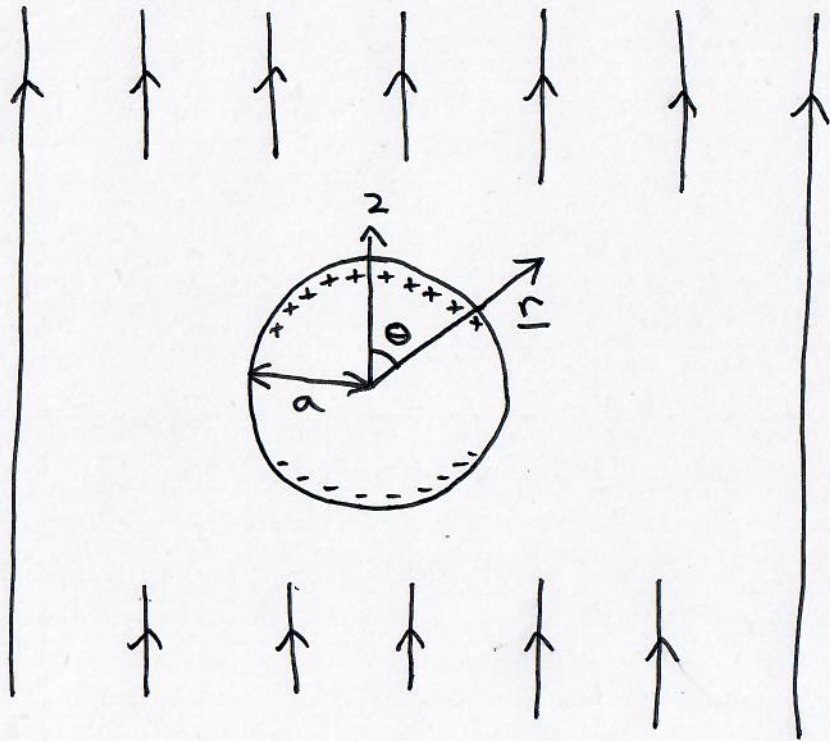


Lecture 4) Example of Laplace's Equation in Spherical Polar Coordinates

$$\underline{\underline{E}} = E_0 \hat{\underline{\underline{z}}}$$



- Potential inside sphere = 0
- E inside sphere = 0

What happens?

- Charge flows
+ve charge at the top of the sphere
-ve charge at the bottom of the sphere
 \Rightarrow Surface of sphere is an equipotential
Let's choose: $V = 0$ at surface of sphere
Centre of sphere at the origin
- At surface of sphere field must be radial
$$\underline{\underline{E}}(r, \theta, \phi) = \underline{\underline{E}}(a, \theta) \hat{\underline{\underline{r}}}$$

\uparrow
ignore because of symmetry in ϕ
- Large distances from the sphere $r \rightarrow \infty$
$$V = -E_0 z = -E_0 r \cos \theta$$
 (in our spherical polar coordinates)

VECTOR DERIVATIVES

Cartesian. $d\mathbf{l} = dx \hat{\mathbf{x}} + dy \hat{\mathbf{y}} + dz \hat{\mathbf{z}}; \quad d\tau = dx dy dz$

Gradient: $\nabla t = \frac{\partial t}{\partial x} \hat{\mathbf{x}} + \frac{\partial t}{\partial y} \hat{\mathbf{y}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z}$

Curl: $\nabla \times \mathbf{v} = \left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) \hat{\mathbf{x}} + \left(\frac{\partial v_x}{\partial z} - \frac{\partial v_z}{\partial x} \right) \hat{\mathbf{y}} + \left(\frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right) \hat{\mathbf{z}}$

Laplacian: $\nabla^2 t = \frac{\partial^2 t}{\partial x^2} + \frac{\partial^2 t}{\partial y^2} + \frac{\partial^2 t}{\partial z^2}$

Spherical. $d\mathbf{l} = dr \hat{\mathbf{r}} + r d\theta \hat{\boldsymbol{\theta}} + r \sin\theta d\phi \hat{\boldsymbol{\phi}}; \quad d\tau = r^2 \sin\theta dr d\theta d\phi$

Gradient: $\nabla t = \frac{\partial t}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial t}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r \sin\theta} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin\theta} \frac{\partial}{\partial \theta} (\sin\theta v_\theta) + \frac{1}{r \sin\theta} \frac{\partial v_\phi}{\partial \phi}$

Curl: $\nabla \times \mathbf{v} = \frac{1}{r \sin\theta} \left[\frac{\partial}{\partial \theta} (\sin\theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{\mathbf{r}}$
 $+ \frac{1}{r} \left[\frac{1}{\sin\theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\boldsymbol{\theta}} + \frac{1}{r} \left[\frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\boldsymbol{\phi}}$

Laplacian: $\nabla^2 t = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial t}{\partial r} \right) + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial t}{\partial \theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2 t}{\partial \phi^2}$

Cylindrical. $d\mathbf{l} = ds \hat{\mathbf{s}} + s d\phi \hat{\boldsymbol{\phi}} + dz \hat{\mathbf{z}}; \quad d\tau = s ds d\phi dz$

Gradient: $\nabla t = \frac{\partial t}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$

Curl: $\nabla \times \mathbf{v} = \left[\frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right] \hat{\mathbf{s}} + \left[\frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right] \hat{\boldsymbol{\phi}} + \frac{1}{s} \left[\frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{\mathbf{z}}$

Laplacian: $\nabla^2 t = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial t}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 t}{\partial \phi^2} + \frac{\partial^2 t}{\partial z^2}$

General Formalism for 2D (r, θ) systems (because of symmetry, in ϕ)

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) = 0 \quad (\text{Eqn 4.1})$$

Employ separation of variables

$$V(r, \theta) = R(r) \Theta(\theta)$$

(Eqn 4.2)

$$\underbrace{\frac{1}{R} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right)}_{\text{constants } l(l+1)} + \underbrace{\frac{1}{\Theta \sin \theta} \frac{d}{d\theta} \left(\sin \theta \frac{d\Theta}{d\theta} \right)}_{-l(l+1)} = 0$$

General solution for equation in \mathbb{R}

$$R(r) = Ar^\lambda + \frac{B}{r^{\lambda+1}}$$

Exercise: verify that this expression is a solution to the equation in \mathbb{R} .

General solution for equation in \mathbb{H}

$$\mathbb{H}(\theta) = P_\lambda(\cos\theta)$$

Exercise: verify that P_0, P_1, P_2, \dots are solutions to the equation in \mathbb{H} .

where $P_\lambda(\cos\theta)$ are the Legendre Polynomials given by:

$$P_0(\cos\theta) = 1$$

$$P_1(\cos\theta) = \cos\theta$$

$$P_2(\cos\theta) = (3\cos^2\theta - 1)/2$$

$$P_\lambda(\cos\theta) = \frac{1}{2^\lambda \lambda!} \left(\frac{d}{d(\cos\theta)} \right)^\lambda (\cos^2\theta - 1)^\lambda$$

General solution for $V(r, \theta)$

$$V(r, \theta) = \sum_{\lambda=0}^{\infty} \left(A_{\lambda} r^{\lambda} + \frac{B_{\lambda}}{r^{\lambda+1}} \right) P_{\lambda}(\cos \theta)$$

$$\int_0^{\pi} V P_{\lambda}(\cos \theta) \sin \theta d\theta = \frac{2}{2\lambda+1} \left(A_{\lambda} r^{\lambda} + \frac{B_{\lambda}}{r^{\lambda+1}} \right)$$

However if V can be expressed as some function $f(\cos \theta)$ it can be simpler to obtain A_{λ} and B_{λ} "by inspection".

Return to our specific problem and apply the boundary conditions

- $r \rightarrow \infty$, $V = -E_0 r \cos \theta$

By comparison with the general solution we can see that

$A_l = 0$ except for the $l=1$ term:

$$\therefore A_1 r \cos \theta = -E_0 r \cos \theta$$

$$V = -E_0 r \cos \theta + \sum \frac{B_l}{r^{l+1}} P_l(\cos \theta)$$

- $r = a$ (surface of sphere), $V = 0$

True for all θ only if $B_l = 0$ except for the $l=1$ term

$$0 = -E_0 a \cos \theta + \frac{B_1}{a^2} \cos \theta \quad \therefore B_1 = E_0 a^3$$

$$\therefore V(r, \theta) = -E_0 \left(r - \frac{a^3}{r^2} \right) \cos \theta$$

External field

Induced charge on sphere

Electric Field

$$\begin{aligned} \underline{E} &= -\nabla V = -\left(\frac{\partial V}{\partial r} \underline{\hat{r}} + \frac{1}{r} \frac{\partial V}{\partial \theta} \underline{\hat{\theta}} \right) \\ &= E_0 \left(1 + \frac{2a^3}{r^3} \right) \cos \theta \underline{\hat{r}} - E_0 \left(1 - \frac{a^3}{r^3} \right) \sin \theta \underline{\hat{\theta}} \end{aligned}$$

At $r=a$:

$$\underline{E} \quad E_0 = 0$$

$$E_r = 3E_0 \cos \theta = \underbrace{\frac{\sigma(\theta)}{\epsilon_0}}$$

from Gauss's theorem

(Since $E=0$ inside the sphere)

$$\sigma(\theta) = 3\epsilon_0 E_0 \cos \theta$$

Note

For systems without symmetry in ϕ solutions in terms of

"Spherical Harmonics" $Y_l^m(\theta, \phi)$

"Extra curricular" reading see, e.g.,

Heald and Marion : Section 3.4

Jackson : Chapter 3.5