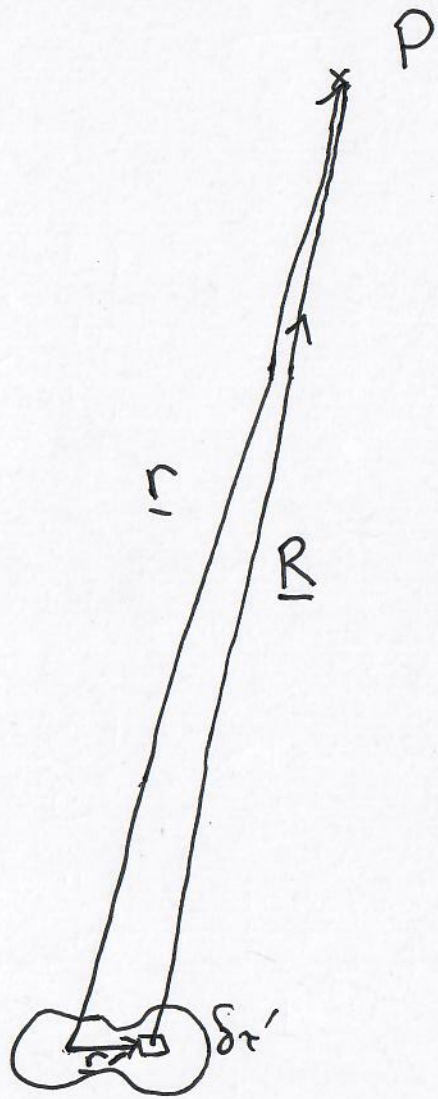


Lecture 5) Multipole Expansions in Electrostatics



$$V(\underline{r}) = \frac{1}{4\pi\epsilon_0} \int_V \frac{\rho(\underline{r}')}{R} d\tau'$$

$$\text{where } \underline{R} = \underline{r} - \underline{r}'$$

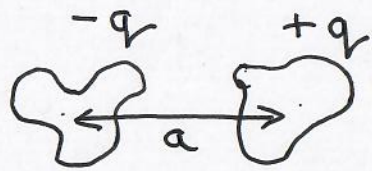
Sometimes it is useful to consider approximate expressions for $V(\underline{r})$ for the case $r \gg r'$

$$\text{Nett charge } Q = \int \rho(\underline{r}') d\tau'$$

$$\text{1st approximation: } V_1 \approx \frac{1}{4\pi\epsilon_0} \frac{Q}{r} \quad \text{for } \underline{R} \approx \underline{r}$$

"monopole" approximation

What if $Q = 0$ overall, but positive and negative charge distributions do not perfectly cancel one another?



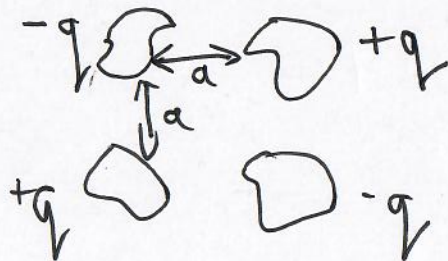
$$r \gg a$$

Can we guess the functional dependence of V on a and r ?

Seems natural to assume that $V \propto \frac{a}{r^2}$

Guess that
$$V = \frac{1}{4\pi\epsilon_0} \frac{qa}{r^2}$$

"dipole" term.



$$V \propto \frac{a^2}{r^3}$$

"quadrupole" term.

Please read Griffiths Introduction to Electrodynamics
page 152-153.

$$\frac{1}{z} \Leftrightarrow \frac{1}{R}$$

Insert series expansion of $\frac{1}{R}$ in terms of powers of $\frac{1}{r}$ into the expression for $V(\underline{r})$ to obtain

$$V(\underline{r}) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{(n+1)}} \int_V (r')^n P_n(\cos\alpha) \rho(\underline{r}') d\tau'$$

N.B. r is a constant as far as the integration is concerned.

Writing out the first few terms, which may be sufficient
for $r \gg r'$

$$V(\underline{r}) = \frac{1}{4\pi\epsilon_0} \left\{ \underbrace{\frac{1}{r} \int \rho(\underline{r}') d\tau'}_{\text{"monopole"}} + \frac{1}{r^2} \int r' \cos\alpha \rho(\underline{r}') d\tau' \right. \\ \left. + \frac{1}{r^3} \int (r')^2 \left[\frac{3}{2} \cos^2\alpha - \frac{1}{2} \right] \rho(\underline{r}') d\tau' \right. \\ \left. + \dots \right.$$

This is called the Multipole expansion for $V(\underline{r})$

Dipole term

$$r' \cos \alpha = \hat{\underline{r}} \cdot \underline{r}'$$

$$V_{\text{dipole}} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\underline{r}} \cdot \int \underline{r}' \rho(\underline{r}') d\tau' = \frac{1}{4\pi\epsilon_0} \frac{\hat{\underline{r}} \cdot \underline{p}}{r^2}$$

where $\underline{p} = \int \underline{r}' \rho(\underline{r}') d\tau'$ is called the "dipole moment"

It depends only on the distribution of charge and not $\underline{r}$!

Consider a point charge q at a position $\underline{a}$ ($a \ll r$)

$$V(\underline{r}) = \frac{q}{4\pi\epsilon_0} \frac{1}{|\underline{r} - \underline{a}|} = \frac{q}{4\pi\epsilon_0} \frac{1}{(r^2 + a^2 - 2ra \cos\alpha)^{1/2}}$$

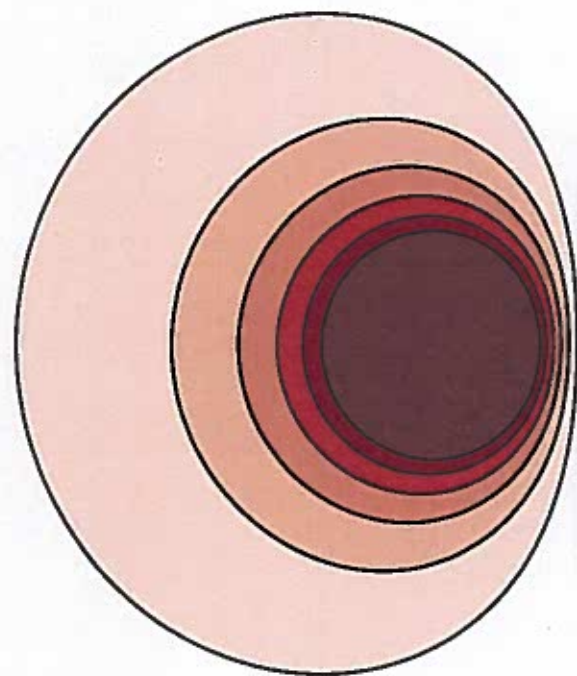
$$\approx \frac{q}{4\pi\epsilon_0} \frac{1}{r} \left(1 - 2 \frac{a}{r} \cos\alpha\right)^{-1/2}$$

$$\approx \frac{q}{4\pi\epsilon_0} \left(\frac{1}{r} + \frac{a \cos\alpha}{r^2} + \dots \right)$$

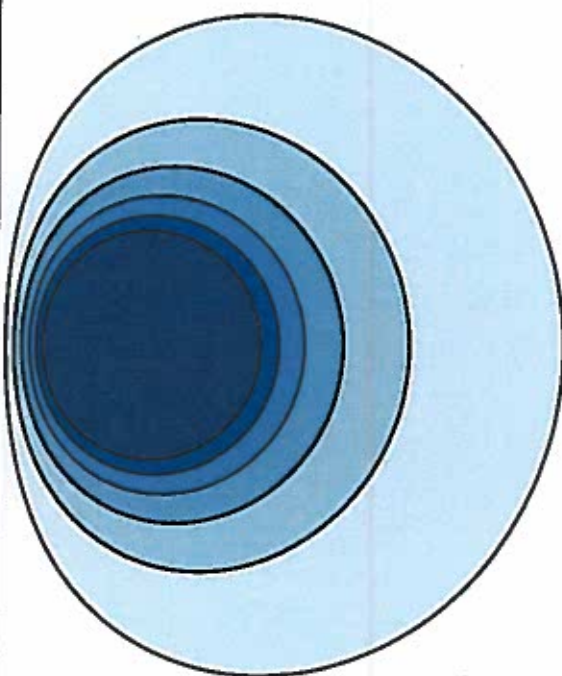
"monopole"
term as if
charge were
at the origin.

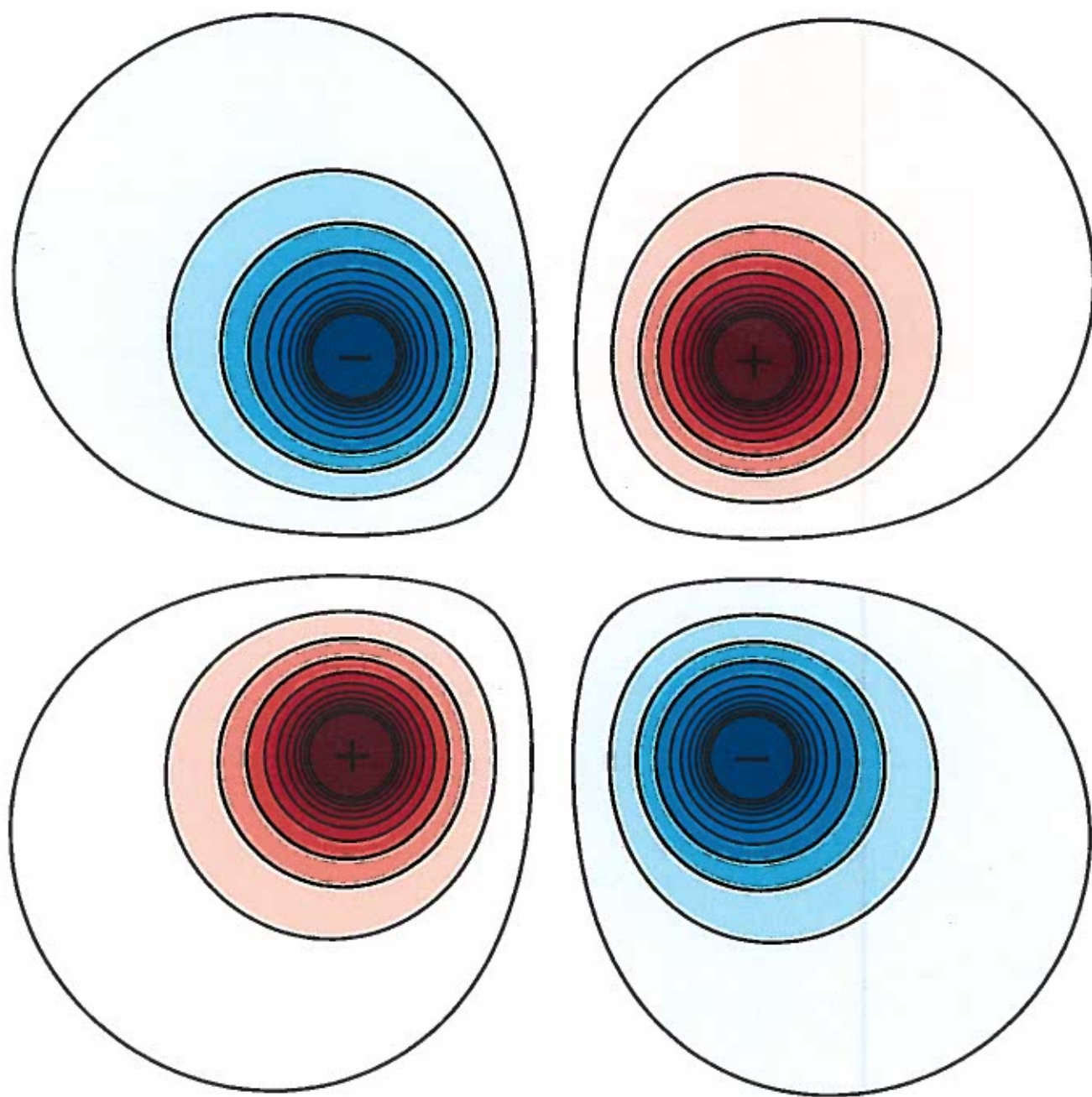
↑
"dipole" and higher terms
appear because charge
is not at the origin

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Suggestions for further reading: Electrostatic Multipoles

Griffiths: "Introduction to Electrodynamics" Chapter 3.4

Heald and Marion: "Classical Electromagnetic Radiation" Chapter 2

Rather more advanced level:

Jackson: "Classical Electrodynamics" Chapter 4 and 5

Suggestions for further reading on magnetostatics for the
next lecture

Griffiths : "Introduction to Electrodynamics" Chapter 5.

Heald and Marion : "Classical Electromagnetic Radiation".

e.g., Sections 1.4, 1.5, 1.7, 2.6 & 2.7.

Rather more advanced level:

Jackson : "Classical Electrodynamics" Chapter 5.