

Lecture 7(c): Multipole expansions in Magnetostatics

Given similarity in form between $V(\underline{r})$ in terms of $\rho(\underline{r}')$
 $\underline{A}(\underline{r})$ $\underline{j}(\underline{r}')$

we can just write down the main results!

$$\underline{A}(\underline{r}) = \frac{\mu_0 I}{4\pi} \int \frac{d\underline{l}'}{R} \quad \text{and} \quad \frac{1}{R} = \frac{1}{r} \sum_{l=0}^{\infty} \left(\frac{r'}{r}\right)^l P_l(\cos\alpha)$$

Consider current flowing in a closed loop



$$\underline{A}(\underline{r}) = \frac{\mu_0 I}{4\pi} \left[\underbrace{\frac{1}{r} \oint d\underline{l}'}_{l=0 \text{ "monopole" term} = 0} + \underbrace{\frac{1}{r^2} \oint r' \cos\alpha d\underline{l}'}_{\text{"dipole" term}} + \underbrace{\frac{1}{r^3} \oint (r')^2 \left[\frac{3\cos^2\alpha - 1}{2} \right] d\underline{l}'}_{\text{"quadrupole" term}} \right]$$

∴ Integral of vector displacement
around a closed path = 0

Notes

FROM LECTURE 6

1) Contribution to $\underline{A}(\underline{r})$ from $\underline{j}(\underline{r}')d\tau'$ is in the same direction as $\underline{j}(\underline{r}')$

2) We can think of eqn 6.1 as effectively three equations that relate components of $A_i(\underline{r})$ to $j_i(\underline{r}')$ for $i=1,2,3$

3) Note the striking similarity in form of:

$$A_i(\underline{r}) = \frac{\mu_0}{4\pi} \int \frac{j_i(\underline{r}')}{R} d\tau' \quad \text{to} \quad V(\underline{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\underline{r}')}{R} d\tau'$$

4) Later in the course we shall show that

$\left. \begin{array}{l} \left(\frac{V}{c}, \underline{A} \right) \\ \text{and } \left(\rho c, \underline{j} \right) \end{array} \right\}$ are 4-vectors and write the equation that relate them in 4-vector form.

Let's take a closer look at the dipole term

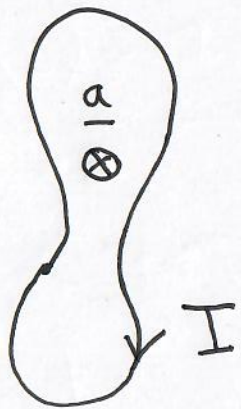
In example sheet 3 you will show that

$$\oint r' \cos \alpha \, d\mathbf{l}' = \oint (\hat{\mathbf{r}} \cdot \mathbf{r}') \, d\mathbf{l}' = -\hat{\mathbf{r}} \times \int_S d\mathbf{a}' = -\hat{\mathbf{r}} \times \mathbf{a}'$$

"vector area"

$$\therefore \underline{A}_{\text{dipole}}(\mathbf{r}) = \frac{\mu_0 I}{4\pi} \frac{\mathbf{a} \times \hat{\mathbf{r}}}{r^2} = \frac{\mu_0}{4\pi} \frac{\mathbf{m} \times \hat{\mathbf{r}}}{r^2}$$

where $\mathbf{m} = I \mathbf{a}$ is the "magnetic dipole moment"

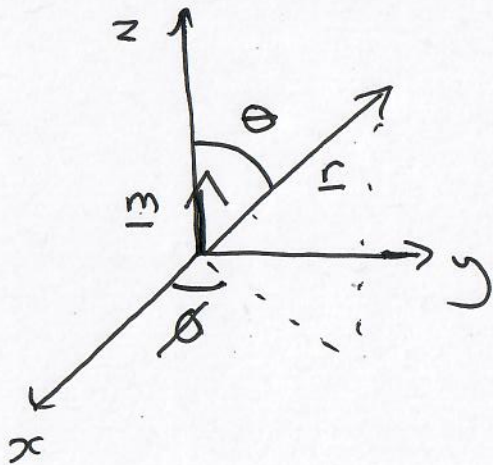


Note: If the loop lies in a plane $|\mathbf{a}|$ is the area enclosed

Direction of \mathbf{a} is given by the circulation of I using the right hand rule.

For example, let's consider $\underline{m} = m \hat{z}$ at the origin.

and work out expressions for \underline{A} and \underline{B}



$$\underline{A}_{\text{dipole}} = \frac{\mu_0 m}{4\pi} \frac{\sin\theta}{r^2} \hat{\phi}$$

$$\begin{aligned} \underline{B} = \nabla \times \underline{A} &= \frac{\mu_0 m}{4\pi} \left[\frac{1}{r \sin\theta} \frac{\partial}{\partial\theta} \left(\sin\theta \frac{\sin\theta}{r^2} \right) \hat{r} - \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\sin\theta}{r^2} \right) \hat{\theta} \right] \\ &= \frac{\mu_0 m}{4\pi} \left[\frac{2 \cos\theta}{r^3} \hat{r} + \frac{\sin\theta}{r^3} \hat{\theta} \right] \end{aligned}$$

VECTOR DERIVATIVES

Cartesian. $d\mathbf{l} = dx \hat{\mathbf{x}} + dy \hat{\mathbf{y}} + dz \hat{\mathbf{z}}; \quad d\tau = dx dy dz$

Gradient: $\nabla t = \frac{\partial t}{\partial x} \hat{\mathbf{x}} + \frac{\partial t}{\partial y} \hat{\mathbf{y}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z}$

Curl: $\nabla \times \mathbf{v} = \left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z} \right) \hat{\mathbf{x}} + \left(\frac{\partial v_x}{\partial z} - \frac{\partial v_z}{\partial x} \right) \hat{\mathbf{y}} + \left(\frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y} \right) \hat{\mathbf{z}}$

Laplacian: $\nabla^2 t = \frac{\partial^2 t}{\partial x^2} + \frac{\partial^2 t}{\partial y^2} + \frac{\partial^2 t}{\partial z^2}$

Spherical. $d\mathbf{l} = dr \hat{\mathbf{r}} + r d\theta \hat{\boldsymbol{\theta}} + r \sin \theta d\phi \hat{\boldsymbol{\phi}}; \quad d\tau = r^2 \sin \theta dr d\theta d\phi$

Gradient: $\nabla t = \frac{\partial t}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial t}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r \sin \theta} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$

Curl: $\nabla \times \mathbf{v} = \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{\mathbf{r}} + \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\boldsymbol{\theta}} + \frac{1}{r} \left[\frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\boldsymbol{\phi}}$

Laplacian: $\nabla^2 t = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial t}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial t}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 t}{\partial \phi^2}$

Cylindrical. $d\mathbf{l} = ds \hat{\mathbf{s}} + s d\phi \hat{\boldsymbol{\phi}} + dz \hat{\mathbf{z}}; \quad d\tau = s ds d\phi dz$

Gradient: $\nabla t = \frac{\partial t}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial t}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$

Divergence: $\nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$

Curl: $\nabla \times \mathbf{v} = \left[\frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right] \hat{\mathbf{s}} + \left[\frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right] \hat{\boldsymbol{\phi}} + \frac{1}{s} \left[\frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{\mathbf{z}}$

Laplacian: $\nabla^2 t = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial t}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 t}{\partial \phi^2} + \frac{\partial^2 t}{\partial z^2}$