

1. In Lecture 14 we used the Lorentz transformation to work out the scalar component of the 4-potential produced by a point charge q in a frame of reference, S , in which it is moving along the x^1 axis with constant speed (in units of c) β :

$$\begin{aligned} A^0 &= \frac{V}{c} = \gamma A'^0 = \frac{q}{4\pi\epsilon_0 c} \gamma \frac{1}{R'} \\ &= \frac{q}{4\pi\epsilon_0 c} \gamma \frac{1}{\left[(\gamma[x^1 - \beta x^0])^2 + (x^2)^2 + (x^3)^2\right]^{\frac{1}{2}}} \end{aligned}$$

Show that the potential A^0 may be expressed also as

$$A^0 = \frac{q}{4\pi\epsilon_0 c R} \left(\frac{1}{1 - \beta^2 \sin^2 \theta} \right)^{\frac{1}{2}},$$

where $\mathbf{R}$ is the vector from the current position $x^1 = \beta x^0, x^2 = 0, x^3 = 0$ of the point charge to the point at which the field is evaluated. The magnitude of R is given by

$$R^2 = [x^1 - \beta x^0]^2 + (x^2)^2 + (x^3)^2.$$

θ is the angle between $\mathbf{R}$ and the direction of motion.

2. (a) Show that $(\mathbf{E} \cdot \mathbf{B})$ is a Lorentz invariant by using the explicit transformation equations for the components of $\mathbf{E}$ and $\mathbf{B}$.
(b) Show that $(E^2 - c^2 B^2)$ is a Lorentz invariant.

Note: It would be interesting to do this in two different ways:

- i. By using the explicit transformation equations for $\mathbf{E}$ and $\mathbf{B}$.
- ii. (*Optional challenge*) The quantity $F_{\mu\nu}F^{\mu\nu}$ is manifestly a Lorentz invariant. (All Lorentz indices in the expression come in contravariant-covariant pairs, which are summed over.) By showing that

$$F_{\mu\nu}F^{\mu\nu} = \frac{-2}{c^2} (E^2 - c^2 B^2)$$

we therefore can prove that $(E^2 - c^2 B^2)$ is a Lorentz invariant.

Notes/hints on calculating $F_{\mu\nu}F^{\mu\nu}$:

- i. Because both 4-vector indices are contracted, $F_{\mu\nu}F^{\mu\nu}$ is a scalar. Because of the behaviour of $F_{\mu\nu}$ and $F^{\mu\nu}$ under Lorentz transformations, $F_{\mu\nu}F^{\mu\nu}$ is guaranteed to be a Lorentz scalar.
- ii. Show that the elements of $F_{\mu\nu}$ are the same as those of $F^{\mu\nu}$ except for the replacement $E_i \rightarrow -E_i$. (The signs of the B_i terms are unchanged.)
- iii. The contraction over the two 4-vector indices produces a sum containing 16 terms. It is, perhaps, simplest just to write out explicitly all 16 terms and add them together.

[I hope you've noticed that you walk past an expression involving $F^{\mu\nu}F_{\mu\nu}$ in the entrance hall to the Schuster Building! (It appears as part of the "Lagrangian Density" of the Standard Model of particle physics¹).

- (c) Suppose that in one inertial frame of reference $\mathbf{B} = 0$ but $\mathbf{E} \neq 0$ (at some point P). Is it possible to find another frame in which the electric field is zero at P ?

¹Subtle advert for the 4th year courses on Field Theory and Gauge Theories ;-)