

1. In Lecture 14 we used the Lorentz transformation to work out the scalar component of the 4-potential produced by a point charge  $q$  in a frame of reference,  $S$ , in which it is moving along the  $x^1$  axis with constant speed (in units of  $c$ )  $\beta$ :

$$\begin{aligned} A^0 &= \frac{V}{c} = \gamma A'^0 = \frac{q}{4\pi\epsilon_0 c} \gamma \frac{1}{R'} \\ &= \frac{q}{4\pi\epsilon_0 c} \gamma \frac{1}{\left[ (\gamma [x^1 - \beta x^0])^2 + (x^2)^2 + (x^3)^2 \right]^{\frac{1}{2}}} \end{aligned}$$

Show that the potential  $A^0$  may be expressed also as

$$A^0 = \frac{q}{4\pi\epsilon_0 c R} \left( \frac{1}{1 - \beta^2 \sin^2 \theta} \right)^{\frac{1}{2}},$$

where  $\mathbf{R}$  is the vector from the current position  $x^1 = \beta x^0, x^2 = 0, x^3 = 0$  of the point charge to the point at which the field is evaluated. The magnitude of  $R$  is given by

$$R^2 = [x^1 - \beta x^0]^2 + (x^2)^2 + (x^3)^2.$$

$\theta$  is the angle between  $\mathbf{R}$  and the direction of motion.

2. (a) Show that  $(\mathbf{E} \cdot \mathbf{B})$  is a Lorentz invariant by using the explicit transformation equations for the components of  $\mathbf{E}$  and  $\mathbf{B}$ .  
 (b) Show that  $(E^2 - c^2 B^2)$  is a Lorentz invariant.

Note: It would be interesting to do this in two different ways:

- i. By using the explicit transformation equations for  $\mathbf{E}$  and  $\mathbf{B}$ .
- ii. (*Optional challenge*) The quantity  $F_{\mu\nu} F^{\mu\nu}$  is manifestly a Lorentz invariant. (All Lorentz indices in the expression come in contravariant-covariant pairs, which are summed over.) By showing that

$$F_{\mu\nu} F^{\mu\nu} = \frac{-2}{c^2} (E^2 - c^2 B^2)$$

we therefore can prove that  $(E^2 - c^2 B^2)$  is a Lorentz invariant.

Notes/hints on calculating  $F_{\mu\nu} F^{\mu\nu}$ :

- i. Because both 4-vector indices are contracted,  $F_{\mu\nu} F^{\mu\nu}$  is a scalar. Because of the behaviour of  $F_{\mu\nu}$  and  $F^{\mu\nu}$  under Lorentz transformations,  $F_{\mu\nu} F^{\mu\nu}$  is guaranteed to be a Lorentz scalar.
- ii. Show that the elements of  $F_{\mu\nu}$  are the same as those of  $F^{\mu\nu}$  except for the replacement  $E_i \rightarrow -E_i$ . (The signs of the  $B_i$  terms are unchanged.)
- iii. The contraction over the two 4-vector indices produces a sum containing 16 terms. It is, perhaps, simplest just to write out explicitly all 16 terms and add them together.

[I hope you've noticed that you walk past an expression involving  $F^{\mu\nu} F_{\mu\nu}$  in the entrance hall to the Schuster Building! (It appears as part of the "Lagrangian Density" of the Standard Model of particle physics<sup>1</sup>).

- (c) Suppose that in one inertial frame of reference  $\mathbf{B} = 0$  but  $\mathbf{E} \neq 0$  (at some point  $P$ ). Is it possible to find another frame in which the electric field is zero at  $P$ ?

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<sup>1</sup>Subtle advert for the 4th year courses on Field Theory and Gauge Theories ;-)