

PHYS30441

Electrodynamics

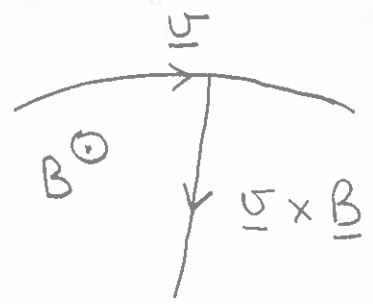
Example Sheet 5

Solutions — Terry Wyatt.

Q1)

(a) For a circular orbit :

$$a = \frac{v^2}{R}$$



In the lectures we showed that the relativistic expression $\frac{dp^\mu}{d\tau} = q F^{\mu\nu} u_\nu$ reduces to

$$\frac{d\underline{p}}{dt} = q \underline{v} \times \underline{B}, \quad \text{since } \gamma d\tau = dt$$

and $u_\nu = (\gamma c, \gamma \underline{v})$

The relativistic expression for momentum is $\underline{p} = \gamma m \underline{v}$

$$\left| \frac{d\underline{p}}{dt} \right| = m \left(v \frac{d\gamma}{dt} + \gamma \frac{dv}{dt} \right) = \gamma m a = \frac{\gamma m v^2}{R} = q v B$$

↑
which is zero
iff $\underline{a} \perp \underline{v}$

$$p = 7 \text{ TeV}/c$$

$$\therefore R = \frac{\gamma m v}{q B} = \frac{p}{q B} = \frac{7 \times 10^{12} \times e}{c} \times \frac{1}{8e}$$

$$\boxed{R \approx 3 \text{ km}}$$

N.B. Can write $a = \frac{q v B}{\gamma m} = \frac{q B p}{(\gamma m)^2}$,

which may be useful for part (b)

Q1)

(b) For synchrotron radiation $\underline{a} \perp \underline{v}$

$$P_{\perp} = \frac{\mu_0 q^2 \gamma^4 a^2}{6\pi c} \quad \text{since } \beta^2 = \left(\frac{a}{c}\right)^2$$

$$= \frac{\mu_0 q^2 \gamma^4}{6\pi c} \times \frac{q^2 B^2}{p^2} \times \left(\frac{p}{\gamma m}\right)^4, \quad \text{since } a^2 = \frac{v^4}{R^2}$$

$$= \frac{\mu_0 q^4 B^2}{6\pi c} \frac{(\gamma mc)^2}{m^4} \quad \text{since } p \approx \gamma mc$$

$$= \frac{\mu_0 q^4 B^2 \gamma^2}{6\pi m^2}$$

Divide by factor
q to get answer
in eV!

γ for proton
with
 $E = 7 \text{ TeV}$
 $mc^2 = 0.94 \text{ GeV}$

$$= \frac{(4\pi \times 10^{-7}) \times (3 \times 10^8)^3 \times (1.6 \times 10^{-19})^3 \times 64 \times \left(\frac{7 \times 10^3}{0.94}\right)^2}{6\pi \times (1.7 \times 10^{-27})^2}$$

$$\approx 10^8 \text{ eV s}^{-1}$$

Q1)

(c) "Circumference" $\approx 2\pi R$

Time for one orbit: $T = \frac{2\pi R}{c}$
(in LHC rest frame)

Energy radiated: $\mathcal{E} = \frac{2\pi R}{c} P_{\perp}$
 $\approx 6.3 \times 10^3 \text{ eV}$

(d) P_{\perp} is a Lorentz invariant quantity

However, time for one orbit in proton rest frame $T' = \frac{T}{\gamma} = \frac{2\pi R}{c\gamma}$

Why?

(i) A clock in the proton rest frame measures "proper time".

or, alternatively,

(ii) The circumference of the LHC is Lorentz contracted in the proton rest frame.

$$\therefore \mathcal{E}' = \frac{\mathcal{E}}{\gamma} \approx 0.8 \text{ eV}.$$

Q2) From the answer to question 1
 $\gamma = \frac{q}{m} \frac{RB}{c}$ for ultra-relativistic particles.

(a) The maximum energy for a given dipole field strength and orbit radius scales as $\frac{\text{charge of the nucleus}}{\text{rest mass}}$ so that for ${}^{208}_{82}\text{Pb}$

we have $E \sim 6.5 \times \frac{82}{208} \sim 2.5 \text{ TeV}$.

(b) Radius of nucleus $\sim \sqrt[3]{\text{number of nucleons}}$

radius for Pb: $r \sim \sqrt[3]{208} \times 10^{-15} \text{ m}$
 $\sim 6 \times 10^{-15} \text{ m}$

(c) Perpendicular to direction of motion

$$E \approx \frac{q\gamma}{4\pi\epsilon_0 r^2} \approx \frac{82 \times 1.6 \times 10^{-19} \times \frac{2.5 \times 10^{12}}{938 \times 10^6}}{4\pi \times 8.85 \times 10^{-12} \times 36 \times 10^{-30}}$$
$$\approx 9 \times 10^{24} \text{ V/m}$$

Q2 (d)

In Lecture 15 we noted that at fixed R'

$$E_{\perp} = E'_{\perp} \quad (\text{equation 15.5})$$

↑
distance from
a point charge
in its rest
frame.

The point at the edge of the nucleus along the direction of motion corresponds to this case.

$$\begin{aligned} \therefore E_{\perp} &\approx \frac{q}{4\pi\epsilon_0 r^2} \approx \frac{82 \times 1.6 \times 10^{-19}}{4\pi \times 8.85 \times 10^{-12} \times 36 \times 10^{-30}} \\ &\approx 3 \times 10^{21} \text{ V/m} \end{aligned}$$

(e) Because distances \perp to the direction of motion are not Lorentz-contracted, the components of $\underline{E} \perp \underline{\beta}$ always gain a factor of γ when we transform from the rest frame (equation 15.6).

In Lecture 15 we noted that at constant R (distance in the frame S , in which the charge is moving) $E = E'/\gamma^2$ for the component of $\underline{E} \parallel \underline{\beta}$!

However, in this problem we were asked for the field at the edge of the nucleus and the distance from centre to edge is Lorentz-contracted